Nek5000 Tutorial

Paul Fischer Aleks Obabko Stefan Kerkemeier Misun Min James Lottes Katie Heisey Shashi Aithal Yulia Peet



Mathematics and Computer Science Division Argonne National Laboratory

Outline

Nek5000 capabilities

- Equations, timestepping, and SEM basics
- Parallel / serial issues that you should understand

Some Resources

Code & Algorithm Tips:

- Nek5000 wiki page (google nek5000)
- www.mcs.anl.gov/~fischer/Nek5000
- Numerical methods & theory:
 - Deville, Fischer, & Mund, Cambridge Univ. Press, 2002



Part I

Nek5000 capabilities

- Gallery
- Brief history
- Equations solved
- Features overview:
 - Spectral element discretization
 - Convergence properties (nek5_svn/examples)
 - Scalability

Applications



Graft

SEM

В

А

С

LDA

Clockwise from upper left:

- Reactor thermal-hydraulics
- Astrophysics
- Combustion
- Oceanography
- Vascular flow modeling





Coarse DNS: Channel Flow at Re_b=13,000

Simulations by J. Ohlsson, KTH, Stockholm



Figure 15: Turbulent channel flow simulations at $Re_{\tau} = 590$ with polynomial order 15 (a resolution of 288 in the homogenous directions and 192 in the wall-normal direction) showing a) mean velocity profile and b) Reynolds stresses together with the turbulent kinetic energy as defined above. In this case only overintegration was used [Case iii)]. Same time step as for the filtered case. Comparison to DNS results (a resolution of 384 in the homogenous directions and 257 in the wall-normal direction) from Moser et al. [19]. DNS data, ----Log Law, -----Nek5000.

Separation in an Asymmetric Diffuser

xh/L

Ohlsson, Schlatter, F., and Henningson,, *JFM* (2010)

Flow separation and recovery DNS at Re=10,000[:] E=127750, N=11, 100 convective time units Comparison w/ exptl results of Cherry et al. 1.00 u=.4U2.00 y 1.0 2.60 Pressure Recovery 2.00 1.00 0.5 1.00 0.4 3.96 $C_{p} 0.3$ Expt 3.00 2.00 0.2 **SEM** 1.00 0.1 0 0.5 1.0 1.5

Axial Velocity



Nek5000 Submission to T-Junction Benchmark

F., Obabko, Tautges, Caceres

- E=62000 spectral elements of order N=7 (n=21 million)
 - Mesh generated by SHARP framework group using CUBIT
 - Loaded into Nek5000 through parallel MOAB interface
 - A group effort at integrated process, under a time constraint...
- Re_D = 40,000 for inlet pipes (exp.: 90,000)
- Subgrid dissipation modeled with low-pass spectral filter
- L_x ~ 25 D (cost is quadratic in L_x)
- 24 hours on 16384 processors of BG/P (850 MHz) ~ 33x slower than uRANS



Instantaneous Temperature Distributions:

Thermal striping leads to thermal fatigue in structural components





Velocity Comparison Downstream of T-junction

- Medium resolution results are in excellent agreement at x=1.6 & 2.6
- Experiment (Re=90K) exhibits more rapid recovery of profile than simulation (Re=40K)



NEA/OECD 2010 Blind T-Junction Benchmark

- Experiment with hot/cold inlets at Re ~ 10⁵
- Inlet velocity and temperature data provided by Vattenfall.
- Results announced 9/15/10 at CFD4NRS meeting in DC
- Nek5000 submission (21 M points) ranked respectively 1st and 6th in temperature and velocity prediction, of 29 entries w/ 1—71 M points.



Pipe Flow at Re=44,000

Motivated by T-junction benchmark:

- Target Re=100,000
 - E=12288, N=11, n~16 million
 - □ ∆t=.0003, CFL ~ 0.35
 - $P_N P_N$
 - Filter-based dissipation: p101,103=3,.001
 - One of many test runs:
 - Dyn. Smag.: good wall shear predictions
 - 1-2 million points
 - No filter

Wu & Moin (2008)

630 m points

- Comparison with:
 - Rudman & Hughes (1999) SEM LES @ Re=36700
 1.2 million points
 - 2nd-order FD DNS Re=44000



Re=44K Pipe Flow Results

- 16 million point filter-based SEM gives perfect Reynolds stress agreement with 630 million point DNS
- **Rigged**...(sort of...): Filter weight chosen to match target WSS.
 - Investigating automatic selection criterion
- 10% variance in friction factor between Princeton & Oregon expts at 44K.
 - DNS closest to combined curve fit.





Rudman & Blackburn (1999) - LES; DenToonder & Nieuwstadt (1997) exp (Re=24,600), Wu & Moin (2008) DNS (Re=44,000)

Resolution Questions: Pipe Flow, Re > 36,700

Plots of w-w*, where w* = low-pass-filtered axial velocity

Illustrates that azimuthal resolution is weak link in current simulations



$$W - W^*$$

 $W - W^*$

| *u* – *u** |

Parallel Scaling: Subassembly 217 Wire-Wrapped Pins

3 million 7th-order spectral elements (n=1.01 billion)
16384–131072 processors of IBM BG/P





www.mcs.anl.gov/~fischer/sem1b

Nek5000 / Star Cross-Channel Velocity Comparison

HEDL geometry $Re_h = 10,500$







Figure 7. Comparison of predicted normal velocity profile for (a) plane group 1 and (b) plane group 3. RANS data is shown in blue and LES data is shown in red.



Figure 8. Comparison of predicted normal velocity profile for (a) plane group 9 and (b) plane group 14. RANS data is shown in blue and LES data is shown in red.

Moving Mesh Examples

peristaltic flow model nek5_svn/examples/peris

2D piston, intake stroke:(15 min. to set up and run)



Moving Mesh Examples

Free surface case

Moving Mesh Examples

Free surface verification:

- agreement with linear theory to 3 significant digits
- long time solution requires filtering



Figure 5: Eigenmodes for free-surface film flow: (left, top) contours of vertical velocity v for unfiltered and (left, bottom) filtered solution at time t = 179.6; (right) error in growth rate vs. t.

Nek5000 Brief History

DNS / LES code for fluid dynamics, heat transfer, MHD, combustion,...

- 100K lines of code: f77 (70K) & C (30K)
- Interfaces w/ VisIt & MOAB/Cubit

Based on high-order spectral element method (Patera '84, Maday & Patera '89)

- Initially Nekton 2.0, first 3D SEM code. (F., Ho, Ronquist, Mavriplis '86-'89)

First commercially-available code for distributed memory computers (marketed by Fluent as Nekton into the mid 90s)

Nek5000 is a highly scalable variant of Nekton

- Gordon Bell Prize in HPC, 4096 processors

(Tufo & F. '99)

- 20% of peak on 262,000 processors of BGP

(Kerkemeier, Parker & F. '10)

Motivation for High-Order

Large problem sizes enabled by peta- and exascale computers allow propagation of small features (size λ) over distances L >> λ .

- Dispersion errors accumulate linearly with time:

~ correct speed – numerical speed * t (for each wavenumber)

 \rightarrow error_{t final} ~ (L / λ) * | numerical dispersion error |

- For fixed final error \mathcal{E}_{f} , require: numerical dispersion error ~ $(\lambda / L)\mathcal{E}_{f}$, << 1
- High-order methods most efficiently deliver small dispersion errors (Kreiss & Oliger 72, Gottlieb et al. 2007)

1D Advection Example

For a complex signal typical of turbulence (-5/3 spectrum), only N >15 case delivers *err* ~ 0.01 after ten revolutions of the signal, out of all cases using 600 points in space.



Spectral Element Convergence: Exponential with N

- Exact Navier-Stokes eigenfunctions with ψ(x, y) comprising:
 cos(mx) cos(ny) cos(mx) sin(ny) sin(mx) cos(ny) sin(mx) sin(ny)
 with m² + n² = λ². [O. Walsh '92]
- $(E, N) = (4^2, 7)$ yields $\epsilon = .01$ after 25 convective time units, $t\bar{U}\lambda/2\pi$.



SEM Excellent transport properties, even for non-smooth solutions



Convection of non-smooth data on a 32x32 grid (K₁ x K₁ spectral elements of order N).

(cf. Gottlieb & Orszag 77)

Strengths of Nek5000

High-order accuracy at low cost

- Extremely rapid (exponential) convergence in space
- 3rd-order accuracy in time

Highly scalable

- Fast scalable multigrid solvers
- Scales to > 290,000 processors with ~10⁴ pts/proc on BGP

Extensively tested

- > 10s of platforms over 25 years
- > 150 journal articles & > 60 users worldwide
- > 400 tests after each build to ensure verified source (more tests to be added)

Solver Performance: Hybrid Schwarz-Multigrid

Magneto-rotational instability

(Obabko, Cattaneo & F.)

- E=140000, N=9 (n = 112 M), P=32768 (BG/L)
- ~ 1.2 sec/step
- ~ 8 iterations / step for U & B
- Key is to have a scalable coarse-grid solver





Scaling to P=262144 Cores

- Production combustion and reactor simulations on ALCF BG/P demonstrate scaling to P=131072 with n/P ~ 5000-10,000 and η ~ .7
- Test problem with 7 billion points scales to P=262144 on Julich BG/P with η ~ .7
 tests 64-bit global addressing for *gs* communication framework



Limitations of Nek5000

No steady-state NS or RANS:

unsteady RANS under development / test – Aithal

Lack of monotonicity for under-resolved simulations

- limits, e.g., LES + combustion
- A high priority for 2011-12

Meshing complex geometries:

- fundamental: meshing always a challenge;

hex-based meshes intrinsically anisotropic

 technical: meshing traditionally not supported as part of advanced modeling development

Mesh Anisotropy

A common refinement scenario (somewhat exaggerated):



cells in the far field

Refinement propagation leads to

- unwanted elements in far-field
- high aspect-ratio cells that are detrimental to iterative solver performance (F. JCP'97)

Some Meshing Options

genbox: unions of tensor-product boxes

prenek: basically 2D + some 3D or 3D via extrusion (n2to3)

Grow your own: 217 pin mesh via matlab; BioMesh



3rd party: CUBIT + MOAB, *TrueGrid*, *Gambit*, *Star CD*

Morphing:



Part 2 (a)

Equations, timestepping, and spectral element formulation

....but first, a bit of code structure.

nek5_svn repository

Key subdirectories in the repo:

- nek5_svn
 - trunk
 - nek makenek script and source files
 - tools several utilities (prenek, genbox, etc.) and scripts
 - examples several case studies

Typical steps to run a case:

- Create a working directory and copy contents of a similar example case to this directory
- Modify case files to suit
- Copy makenek from nek and type makenek <case>
- Run job using a script (tools/scripts) and analyze results (postx/Vislt)

nek5_svn repository

nek5_svn	nek5 svn
3rd_party	:
branches	:
examples	` trunk
axi	nek
benard	i I :
	source files
eddy	ii :
fs_2	` tools
fs_hydro	amg_matlab
kovasznay	avg
lowMach_test	genbox
	genmap
	makefile
	maketools
rayleigh	n2to3
shear4	nekmerge
timing	postnek
turbChannel	prenek
turbJet	reatore2
` vortex	` scripts
tags	
tests	

`-- trunk

Base Nek5000 Case Files

SIZE – an f77 include file that determines

- spatial dimension (Idim =2 or 3)
- approximation order (lx1,lx2,lx3,lxd) N := lx1-1
- upper bound on number of elements per processor: lelt
- upper bound on total number of elements, lelg
- <case>.rea a file specifying
 - job control parameters (viscosity, dt, Nsteps, integrator, etc.)
 - geometry element vertex and curvature information
 - boundary condition types
 - restart conditions

<case>.usr – f77 source file specifying

- initial and boundary conditions
- variable properties
- forcing and volumetric heating
- geometry morphing
- data analysis options: min/max, runtime average, rms, etc.

Snapshot of SIZE

С

С

C C

C

```
parameter (ldim=2)
parameter (lx1=14,ly1=lx1,lz1=1,lelt=80,lelv=lelt)
parameter (lxd=20,lyd=lxd,lzd=1)
parameter (lelx=1,lely=1,lelz=1)
NOTE: for IBM BLUE GENE LX1, LXD has to be an even number (double hummer)
parameter (ldimt= 1)
                        ! upper limit for passive scalars + T
parameter (lp = 64) ! upper limit for number of CPUs
parameter (lelg = 5000) ! upper limit for total number of elements
parameter (lzl=3 + 2*(ldim-3))
parameter (lx2=lx1-0)
parameter (ly2=ly1-0)
parameter (lz2=lz1)
parameter (1x3=1x2)
parameter (1y3=1y2)
parameter (lz3=lz2)
```

Snapshots of .rea file

Parameters section

2 DIMENSIONAL RUN 118 PARAMETERS FOLLOW 1.00000 P001: DENSITY -40.0000P002: VISCOS 0.000000E+00 P003: 0.000000E+00 P004: 0.000000E+00 P005: 0.000000E+00 P006: 1.00000 P007: RHOCP 1.00000 P008: CONDUCT 0.000000E+00 P009: 0.000000E+00 P010: FINTIME 2000.00 P011: NSTEPS -0.100000E-02 P012: DT 0.000000E+00 P013: IOCOMM 0.000000E+00 P014: IOTIME 0.000000E+00 P015: IOSTEP 0.000000E+00 P016: PSSOLVER: 0=default P017: 1.00000 0.500000E-01 P018: GRID < 0 --> # cells -1.00000P019: INTYPE 4.00000 P020: NORDER 0.100000E-05 P021: DIVERGENCE 0.100000E-09 P022: HELMHOLTZ 0.000000E+00 P023: NPSCAL 0.000000E+00 P024: TOLREL 0.000000E+00 P025: TOLABS 2.00000 P026: COURANT/NTAU 3.00000 P027: TORDER

Geometry and boundary conditions

	EI	LEME	NT 5	5[1] GR	OUP 0			
-(.5000	0000	0.0000	00E+00	0.00000	00E+00 -0.	5000000		
0	.5000	0000	0.50000	000	1.0000	00 1	.000000		
	EI	LEME	NT (5[1] GR	OUP 0			
0	0.000	0000	E+00 1.0000	000	1.0000	00 0.	0000000	E+00	
(.5000	0000	0.50000	000	1.0000	00 1	.000000		
	EI	LEME	NT	7[1] GR	OUP 0			
-0	.5000	0000	0.0000	00E+00	0.00000	00E+00 -0.	5000000		
	1.000	0000	1.0000	000	1.5000	00 1	.500000		
	EI	LEME	NT 8	3[1] GR	OUP 0			
(0.000	0000	E+00 1.0000	000	1.0000	00 0.	0000000	E+00	
	1.000	0000	1.0000	000	1.5000	00 1	.500000		
**** CURVED SIDE DATA ****									
		0	Curved sides	follow :	IEDGE,IE	L, CURVE(I)	,I=1,5,	CCURVE	
1	****	BOU	NDARY CONDITI	ONS ***	* *				
1	****	FLU	ID BOUNDARY	CONDIT	IONS ***	**			
Ρ	1	1	7.00000	3.00	000	0.00000E	+00 0.	000000E+00	
Е	1	2	2.00000	4.00	000	0.00000E	+00 0.	000000E+00	
Е	1	3	3.00000	1.00	000	0.00000E	+00 0.	000000E+00	
v	1	4	0.00000E+00	0.000	000E+00	0.00000E	+00 0.	000000E+00	
Ρ	2	1	8.00000	3.00	000	0.00000E	+00 0.	000000E+00	
v	2	2	0.00000E+00	0.000	000E+00	0.00000E	+00 0.	000000E+00	
Е	2	3	4.00000	1.00	000	0.00000E	+00 0.	000000E+00	
Е	2	4	1.00000	2.00	000	0.00000E	+00 0.	000000E+00	
Е	3	1	1.00000	3.00	000	0.00000E	+00 0.	000000E+00	
Е	3	2	4.00000	4.00	000	0.00000E	+00 0.	000000E+00	
Е	3	3	5.00000	1.00	000	0.00000E	+00 0.	000000E+00	
v	3	4	0.000000E+00	0.000	000E+00	0.00000E	+00 0.	000000E+00	
Е	4	1	2.00000	3.00	000	0.00000E	+00 0.	000000E+00	
v	4	2	0.00000E+00	0.000	000E+00	0.00000E	+00 0.	000000E+00	
E	4	3	6.00000	1.00	000	0.000000E	+00 0.	000000E+00	
Snapshot of .usr file

```
_____
c-----
    subroutine userf (ix,iy,iz,eg)
    include 'SIZE'
    include 'TOTAL'
    include 'NEKUSE'
    integer e, f, eq
   e = gllel(eg)
С
    Note: this is an acceleration term, NOT a force!
С
    Thus, ffx will subsequently be multiplied by rho(x,t)
C
    ffx = 0.0
    ffv = 0.0
    ffz = 0.0
    return
    end
c_____
    subroutine userchk
    include 'SIZE'
    include 'TOTAL'
    return
    end
c-----
    subroutine userbc (ix, iy, iz, iside, ieg)
    include 'SIZE'
    include 'TOTAL'
    include 'NEKUSE'
    ux=0.0
    uy=0.0
    uz=0.0
    temp=0.0
    return
    end
                _____
C-
    subroutine useric (ix,iy,iz,ieq)
```

Derived Nek5000 Case Files

- <case>.re2 binary file specifying
 - geometry element vertex and curvature information
 - boundary condition types

This file is not requisite for small problems but important for element counts E > -10,000

<case>.map – ascii file derived from .rea/.re2 files specifying

- mesh interconnect topology
- element-to-processor map

This file is needed for each run and is generated by running the "genmap" tool (once, for a given .rea file).

amg...dat – binary files derived from .rea/.re2 files specifying

algebraic multigrid coarse-grid solver parameters
 These files are needed only for large processor counts (P > 10,000)
 and element counts (E > 50,000).

Part 2 (b)

Equations, timestepping, and spectral element formulation

Outline

Nek5000 capabilities

- Equations, timestepping, and SEM basics
- Parallel / serial issues that you should understand

Equation Sets (2D/3D)

Incompressible Navier-Stokes plus energy equation

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \nabla \cdot \mu \left(\nabla \mathbf{u} + \nabla^T \mathbf{u} \right) + \mathbf{f}$$
$$\nabla \cdot \mathbf{u} = 0$$
$$\rho C_p \left(\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T \right) = \nabla \cdot k \nabla T + q'''$$

plus additional passive scalars:

$$\rho C_{p_i} \left(\frac{\partial T_i}{\partial t} + \mathbf{u} \cdot \nabla T_i \right) = \nabla \cdot k_i \nabla T_i + q_i''', \quad i = 3, ..., n_{flds}$$

Also supports incompressible MHD, low Mach-number hydro, free-surface, and conjugate heat transfer formulations.

Steady State Equations

Steady Stokes (plus boundary conditions):

$$-\nabla \cdot \mu(\mathbf{x}) \left(\nabla \mathbf{u} + \nabla^T \mathbf{u} \right) + \nabla p = \mathbf{f}(\mathbf{x})$$
$$\nabla \cdot \mathbf{u} = 0$$

Steady conduction (plus boundary conditions):

$$-\nabla \cdot k(\mathbf{x}) \nabla T + \lambda(\mathbf{x}) T = q'''(\mathbf{x}), \ \lambda \ge 0$$

Constant Property Equation Set

Incompressible Navier-Stokes + energy equation

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla p + \mu \nabla^2 \mathbf{u} + \mathbf{f}$$
$$\nabla \cdot \mathbf{u} = 0$$
$$\rho C_p \left(\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T \right) = k \nabla^2 T + q'''$$

In Nek parlance, material properties specified in .rea file as:

dimensionalnondimensional (convective time scale) $\bullet p1 = \rho$ $\bullet p1 = 1$ $\bullet p2 = \mu$ $\bullet p2 = 1/Re \text{ (or } -Re)$ $\bullet p7 = \rho C_p$ $\bullet p7 = 1$ $\bullet p8 = k$ $\bullet p8 = 1/Pe \text{ (or } -Pe)$

or as variable properties in f77 routine uservp() (.usr file)

Nek provides a scalable framework to advance these equations with user-defined properties. LES & RANS can be incorporated in this framework. (See /examples.)

Incompressible MHD

$$\frac{\partial \mathbf{u}}{\partial t} - \frac{1}{Re} \nabla^2 \mathbf{u} + \nabla p = \mathbf{B} \cdot \nabla \mathbf{B} - \mathbf{u} \cdot \nabla \mathbf{u},$$
$$\nabla \cdot \mathbf{u} = \mathbf{0}$$
$$\frac{\partial \mathbf{B}}{\partial t} - \frac{1}{Rm} \nabla^2 \mathbf{B} + \nabla q = \mathbf{B} \cdot \nabla \mathbf{u} - \mathbf{u} \cdot \nabla \mathbf{B},$$
$$\nabla \cdot \mathbf{B} = \mathbf{0}$$

- plus appropriate boundary conditions on \boldsymbol{u} and \boldsymbol{B}
- **Typically**, Re >> Rm >> 1
- Semi-implicit formulation yields independent Stokes problems for u and B

Timestepping

Navier-Stokes Time Advancement

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u}$$
$$\nabla \cdot \mathbf{u} = 0$$

Nonlinear term: explicit via BDFk/EXTk or characteristics (Pironneau '82)

- Linear Stokes problem: pressure/viscous decoupling:
 - 3 Helmholtz solves for velocity ("easy" w/ Jacobi-preconditioned CG)
 - (consistent) Poisson equation for pressure (computationally dominant)
- Filter velocity (& temperature), if required

(Order is slightly different for $P_N - P_N$.)

Timestepping Design

Implicit:

- symmetric and (generally) linear terms,
- fixed flow rate conditions

Explicit:

- nonlinear, nonsymmetric terms,
- user-provided rhs terms, including
 - Boussinesq and Coriolis forcing

Rationale:

- div $\mathbf{u} = 0$ constraint is fastest timescale
- Viscous terms: explicit treatment of 2nd-order derivatives $\rightarrow \Delta t \sim O(\Delta x^2)$
- Convective terms require only $\Delta t \sim O(\Delta x)$
- For high Re, temporal-spatial accuracy dictates $\Delta t \sim O(\Delta x)$
- Linear symmetric is "easy" nonlinear nonsymmetric is "hard"

BDF2/EXT2 Example

Consider the convection-diffusion equation,

$$\frac{\partial u}{\partial t} + \mathbf{c} \cdot \nabla u = \nu \nabla^2 u.$$

Discretize in space:

$$B\frac{d\underline{u}}{dt} + C\underline{u} = -\nu A\underline{u}, \qquad (A \text{ is SPD})$$

Evaluate each term at t^n according to convenience:

$$B\frac{d\underline{u}}{dt}\Big|_{t^n} = B\frac{3\underline{u}^n - 4\underline{u}^{n-1} + \underline{u}^{n-2}}{2\Delta t} + O(\Delta t^2)$$

$$C\underline{u}\Big|_{t^n} = 2C\underline{u}^{n-1} - C\underline{u}^{n-2} + O(\Delta t^2)$$

 $\nu A \underline{u} \Big|_{t^n} = \nu A \underline{u}^n$

BDFk/EXTk

BDF3/EXT3 is essentially the same as BDF2/EXT2

- $O(\Delta t^3)$ accuracy
- essentially same cost
- accessed by setting Torder=3 (2 or 1) in .rea file
- For convection-diffusion and Navier-Stokes, the "EXTK" part of the timestepper implies a CFL (Courant-Friedrichs-Lewy) constraint

$$\max_{\mathbf{x}\in\Omega}\frac{|\mathbf{u}|\Delta t}{\Delta x} \approx 0.5$$

- For the spectral element method, $\Delta x \sim N^{-2}$, which is restrictive.
 - We therefore often use a characteristics-based timestepper.
 (IFCHAR = T in the .rea file)

Characteristics Timestepping

Apply BDFk to material derivative, e.g., for k=2:

$$\frac{Du}{Dt} := \frac{\partial u}{\partial t} + \mathbf{c} \cdot \nabla u$$
$$= \frac{3u^n - 4\tilde{u}^{n-1} + \tilde{u}^{n-2}}{2\Delta t} + O(\Delta t^2)$$

Amounts to finite-differencing along the characteristic leading into x_j



Characteristics Timestepping

Don't need <u>position</u> (e.g., Xⁿ⁻¹) of characteristic departure point, only the <u>value</u> of uⁿ⁻¹(x) at these points.

These values satisfy the pure hyperbolic problem:

$$\begin{aligned} \frac{\partial \tilde{u}}{\partial s} + \mathbf{c} \cdot \nabla \tilde{u} &= 0, \qquad s \in [t^{n-1}, t^n] \\ \tilde{u}(\mathbf{x}, t^{n-1}) &:= u^{n-1}(\mathbf{x}), \end{aligned}$$

which is solved via explicit timestepping with $\Delta s \sim \Delta t_{CFL}$

Spatial Discretization

Spectral Element Method

- Variational method, similar to FEM, using *GL* quadrature.
- Domain partitioned into E high-order quadrilateral (or hexahedral) elements (decomposition may be nonconforming - *localized refinement*)
- Trial and test functions represented as N th-order tensor-product polynomials within each element. ($N \sim 4 15$, typ.)
- **E** N^3 gridpoints in 3D, EN^2 gridpoints in 2D.
- Converges *exponentially fast* with *N* for smooth solutions.



Spectral Element Method: Poisson Example

- The SEM is a weighted residual method.
- Consider Poission eqn:

$$-\nabla^2 u = f, \ u|_{\partial} \Omega = 0$$

• Postulate a representation of the solution, e.g.,

$$u(\mathbf{x}) = \sum_{j=1}^{n} u_j \phi_j(\mathbf{x}),$$

and insist that the residual $r(\mathbf{x}) := f + \nabla^2 u$ be orthogonal to all functions in the approximation space:

$$\int_{\Omega} v(f + \nabla^2 u) d\mathbf{x} = 0, \text{ for } v = \phi_i, \ i = 1, \dots, n$$

Spectral Element Method: Poisson Example

Integrate 2nd-order term by parts:

$$-\int_{\Omega} v \nabla^2 u d\mathbf{x} = \int_{\Omega} \nabla v \cdot \nabla u d\mathbf{x} - \int_{\partial \Omega} v \nabla u \cdot \hat{\mathbf{n}} dA$$

- Surface integral vanishes because of boundary conditions.
- Rearranging and inserting basis functions yields:

$$\sum_{j=1}^{n} a_{ij} u_j = \sum_j b_{ij} f_j \iff A \underline{u} = B \underline{f}$$
$$a_{ij} := \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j d\mathbf{x} \qquad \text{stiffness matrix}$$
$$b_{ij} := \int_{\Omega} \phi_i \phi_j d\mathbf{x} \qquad \text{mass matrix}$$
$$f(\mathbf{x}) := \sum_{j=1}^{n} f_j \phi_j(\mathbf{x})$$

SEM Function Representation

Key point is that there is a continuous representation of all variables:

$$u(\mathbf{x}) = \sum_{j=1}^{n} u_j \phi_j(\mathbf{x})$$

Since $\phi_i(\mathbf{x})$ is known a priori, we know how to differentiate and integrate.

Moreover, choose ϕ_i s to be computationally convenient

1D Basis Functions

$$u(x) := \sum_{i=0}^{N} u_i h_i(x), \quad h_i(x) \in \mathbb{P}_N$$

Stable high-order basis for Nth-order polynomial approximation space:

poor choices:

 $h_i(x) = x^i$

 $h_i(x) = \text{Lagragian interpolant on uniform points}, x_i = i \cdot \Delta x$

good choices:

 $h_i(x) = L_i(x)$ (any orthogonal polynomial) $h_i(x) = Lagragian$ interpolant on Gauss points

$$h_i(x)$$

Condition Number of 1D Stiffness Matrix

GLL Nodal Basis \rightarrow good conditioning, minimal round-off error



Integration (projection) in 1D

Replace integrals with Gauss-Lobatto-Legendre quadrature:

$$\int_{\Omega} \nabla v \cdot \nabla u \, dV = \int_{\Omega} v f \, dV$$

with

$$(\nabla v, \nabla u)_N = (v, f)_N$$

where

$$(f,g)_N := \sum_{k=0}^N \rho_k f(\xi_k) g(\xi_k),$$

 $\{\xi_k\}$ = Gauss-Lobbato-Legendre quadrature points

 $\{\rho_k\}$ = Gauss-Lobbato-Legendre quadrature weights

 Yields a diagonal mass matrix; preserves spectral accuracy. (However, beware stability issues....)

Extension to 2D

Nodal bases on the Gauss-Lobatto-Legendre points:





Matrix-Matrix Based Derivative Evaluation

Local tensor-product form (2D),

$$u(r,s) = \sum_{i=0}^{N} \sum_{j=0}^{N} u_{ij}h_i(r)h_j(s), \quad h_i(\xi_p) = \delta_{ip}, \ h_i \in \mathbb{P}_N$$

$$h_i(r)$$

allows derivatives to be evaluated as matrix-matrix products:

$$\frac{\partial u}{\partial r}\Big|_{\xi_i,\xi_j} = \sum_{p=0}^N u_{pj} \frac{dh_p}{dr}\Big|_{\xi_i} = \sum_p \underbrace{\hat{D}_{ip} u_{pj}}_{\text{mxm}} =: D_r \underline{u}$$

■ In 3D, ~90% of all ops in mxm – high optimization important.

Mapped Geometries



Geometry takes same form as solution,

$$\mathbf{x}(r,s) = \sum_{i=0}^{N} \sum_{j=0}^{N} \mathbf{x}_{ij} h_i(r) h_j(s),$$

with $\mathbf{x} := (x, y)$, from which

$$\left. \frac{\partial x}{\partial r} \right|_{\xi_i,\xi_j} = \sum_p D_{ip} x_{pj}$$

Given $\frac{\partial x_i}{\partial r_j}$, we can find $\frac{\partial r_i}{\partial x_j}$, and thus use the chain rule, e.g., $\frac{\partial u}{\partial x} = \frac{\partial u}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial u}{\partial s} \frac{\partial s}{\partial x}.$

In Nek: "call gradm1(ux,uy,uz,u)" assuming ux,uy,uz,u properly declared.

Notes about Mapped Elements

- Best to use affine (i.e., linear) transformations in order to preserve underlying GLL spacing for stability and accurate quadrature.
- Avoid singular corners ~180° or ~0°
- Avoid high-aspect-ratio cells, if possible



Multidimensional Integration

Given that we have Lagrangian interpolants based on GLL quadrature points, we have

$$\sum_{n=1}^{n} vud\mathbf{x} \approx \sum_{k=1}^{n} \rho_k v(\xi_k) u(\xi_k)$$
$$= \sum_k \rho_k \left(\sum_{i=1}^{n} v_i \phi_i(\xi_k) \right) \left(\sum_{j=1}^{n} u_j \phi_j(\xi_k) \right)$$
$$= \sum_k \rho_k \left(\sum_{i=1}^{n} v_i \delta_{ik} \right) \left(\sum_{j=1}^{n} u_j \delta_{jk} \right)$$
$$= \sum_k \rho_k v_k u_k = \underline{v}^T B \underline{u}, \ b_{ij} := \delta_{ij} \rho_i.$$

In particular,

$$\int_{\Omega} u d\mathbf{x} = \sum_{k} \rho_{k} u_{k} = \underline{b}^{T} \underline{u}, \ b_{i} := \rho_{i}.$$

In Nek, this vector reduction is implemented as: alpha = glsc2(u,bm1,n)

Navier-Stokes Discretization Options

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u}$$
$$\nabla \cdot \mathbf{u} = 0$$

- Imposition of the constraint div u = 0 is a major difficulty in solving the incompressible Navier-Stokes equations, both from theoretical and implementation perspectives.
- Was not well-understood till the mid-80s (give, or take...).
- The fundamental difficulty is that the discrete operators do not commute, except under special circumstances (e.g., Fourier bases).
- Nek supports two distinct approaches:
 - Option 1 ($P_N P_{N-2}$):
 - discretize in space using compatible approximation spaces
 - solve coupled system for pressure/velocity
 - Option 2 ($P_N P_N$, or *splitting*):
 - discretize in time first
 - take continuous divergence of momentum equation to arrive at a Poisson equation for pressure, with special boundary conditions

P_N - **P**_{N-2} Spectral Element Method for Navier-Stokes (MP 89)

WRT: Find
$$\mathbf{u} \in X^N$$
, $p \in Y^N$ such that:

$$\frac{1}{Re} (\nabla \mathbf{u}, \nabla \mathbf{v})_{GL} + \frac{1}{\Delta t} (\mathbf{u}, \mathbf{v})_{GL} - (p, \nabla \cdot \mathbf{v})_G = (\mathbf{f}, \mathbf{v})_{GL} \quad \forall \mathbf{v} \in X^N \subset H^1$$

$$- (q, \nabla \cdot \mathbf{u})_G = 0 \qquad \forall q \in Y^N \subset L^2$$

Velocity, **u** in P_N , continuous Pressure, p in P_{N-2} , discontinuous



Gauss-Lobatto Legendre points (velocity)



Consistent Splitting for Unsteady Stokes (MPR 90, Blair-Perot 93, Couzy 95)

$$\begin{bmatrix} \mathbf{H} & \mathbf{D}^T \\ -\mathbf{D} & 0 \end{bmatrix} \begin{pmatrix} \underline{\mathbf{u}}^n \\ \underline{p}^n - \underline{p}^{n-1} \end{pmatrix} = \begin{pmatrix} \mathbf{B}\underline{\mathbf{f}} + \mathbf{D}^T \underline{p}^{n-1} \\ \underline{f}_p \end{pmatrix}$$

$$\begin{bmatrix} \mathbf{H} & -\frac{\Delta t}{\beta_0} \mathbf{H} \mathbf{B}^{-1} \mathbf{D}^T \\ \mathbf{0} & E \end{bmatrix} \begin{pmatrix} \underline{\mathbf{u}}^n \\ \underline{p}^n - \underline{p}^{n-1} \end{pmatrix} = \begin{pmatrix} \mathbf{B} \underline{\mathbf{f}} + \mathbf{D}^T \underline{p}^{n-1} \\ \underline{g} \end{pmatrix} + \begin{pmatrix} \underline{\mathbf{r}} \\ \underline{0} \end{pmatrix} .$$

$$E := \frac{\Delta t}{\beta_0} \mathbf{D} \mathbf{B}^{-1} \mathbf{D}^T \checkmark \mathbf{r} = O(\Delta t^2)$$

- *E* consistent Poisson operator for pressure, SPD
 - boundary conditions applied in velocity space
 - most compute-intensive phase

Comparison of $P_N - P_{N-2}$ and $P_N - P_N$ Options in Nek

		<u>P</u> _N - P _{N-2}	$\underline{P}_{\underline{N}} - \underline{P}_{\underline{N}}$
_	SIZE:	lx2=lx1-2	lx2=lx1
_	pressure:	discontinuous	continuous
_	solver:	$E = DB^{-1}D^{T}$	A (std. Laplacian)
_	preconditioner:	SEMG	Schwarz (but to be upgraded)
_	free-surface	Yes	No
_	ALE	Yes	No
_	low Mach	No	Yes
_	LES	OK	Better
_	low Re	Better	OK
_	var. prop.	Implicit (stress formulation)	semi-implicit
_	spectrally accurate	Yes	Yes

- Nek will ensure that the problem type is compatible with the discretization choice.
- For most cases, speed is determined by the pressure solve, which addresses the fastest timescales in the system (the acoustic waves).
 - For $P_N P_{N-2}$, the solver has been highly optimized over the last 15 years.
 - The P_N P_N version was developed by the ETH group (Tomboulides, Frouzakis, Kerkemeier) for low Mach-number combustion and has only recently been folded into the production Nek5000 code.

Navier-Stokes Boundary Conditions

A few key boundary conditions are listed below.

cbc	name	condition
V	velocity	specified in .usr
V	velocity	specified in .rea
W	wall	$\mathbf{u} = 0$
0	outflow	$\frac{\partial \mathbf{u}}{\partial \mathbf{n}} = 0, \ p = 0$
SYM	symmetry	$\frac{\partial u_t}{\partial \mathbf{n}} = 0, \ u_n = 0$
Ρ	periodic	u(x) = u(x + L)

There are many more, particularly for moving walls, free surface, etc.

Special conditions include:

- Recycling boundary conditions (special form of "v")
- Accelerated outflow to avoid incoming characteristics

Thermal Boundary Conditions

A few key boundary conditions are listed below.

cbc	name	condition
t	temperature	specified in .usr
т	temperature	specified in .rea
Ι	insulated	$\frac{\partial T}{\partial \mathbf{n}} = 0$
f	flux	$k \frac{\partial T}{\partial \mathbf{n}} = f$
С	Newton cooling	$k \frac{\partial T}{\partial \mathbf{n}} = h(T - T_{\infty})$
0	outflow	$\frac{\partial T}{\partial \mathbf{n}} = 0$
Ρ	periodic	$T(\mathbf{x}) = T(\mathbf{x} + \mathbf{L})$

Filtering, Dealiasing, and LES

LES, or more properly, subgrid scale (sgs) stress models invariably rely on dissipation to remove energy that otherwise piles up at the gridscale. WHY?

Stabilizing High-Order Methods

In the absence of eddy viscosity, some type of stabilization is generally required at high Reynolds numbers.

Some options:

- high-order upwinding (e.g., DG, WENO)
- bubble functions
- spectrally vanishing viscosity
- filtering
- dealiasing
Filter-Based Stabilization

(Gottlieb et al., Don et al., Vandeven, Boyd, ...)

At end of each time step:

- Interpolate u onto GLL points for P_{N-1}
- Interpolate back to GLL points for P_N

 $F_1(u) = I_{N-1} u$



- Results are smoother with linear combination:
Mullen 01)

$$F_{\alpha}(u) = (1-\alpha) u + \alpha I_{N-1} u$$
 ($\alpha \sim 0.05 - 0.2$)

- Post-processing no change to existing solvers
- Preserves interelement continuity and spectral accuracy
- Equivalent to multiplying by $(1-\alpha)$ the *N* th coefficient in the expansion

 $- u(x) = \sum u_k \phi_k(x) \rightarrow u^*(x) = \sum \sigma_k u_k \phi_k(x), \quad \sigma_k = 1, \quad \sigma_N = (1 - \alpha)$ $\Box \phi_k(x) := L_k(x) - L_{k-2}(x) \quad (Boyd 98)$

Numerical Stability Test: Shear Layer Roll-Up (Bell et al. JCP 89, Brown & Minion, JCP 95, F. & Mullen, CRAS 2001)



Figure 1: Vorticity for different (K, N) pairings: (a-d) $\rho = 30$, $Re = 10^5$, contours from -70 to 70 by 140/15; (e-f) $\rho = 100$, Re = 40,000, contours from -36 to 36 by 72/13. (cf. Fig. 3c in [4]).

Error in Predicted Growth Rate for Orr-Sommerfeld Problem at Re=7500

(Malik & Zang 84)

Spatial and Temporal Convergence (F. & Mullen, 01)									
	$\Delta t = 0.003125$		N = 17	2nd Order		3rd Order			
N	lpha=0.0	lpha=0.2	Δt	lpha=0.0	lpha=0.2	lpha=0.0	lpha=0.2		
7	0.23641	0.27450	0.20000	0.12621	0.12621	171.370	0.02066		
9	0.00173	0.11929	0.10000	0.03465	0.03465	0.00267	0.00268		
11	0.00455	0.01114	0.05000	0.00910	0.00911	161.134	0.00040		
13	0.00004	0.00074	0.02500	0.00238	0.00238	1.04463	0.00012		
15	0.00010	0.00017	0.01250	0.00065	0.00066	0.00008	0.00008		



Base velocity profile and perturbation streamlines

Filtering permits Re_{*899} > 700 for transitional boundary layer calculations*</sub>



Why Does Filtering Work ? (Or, Why Do the Unfiltered Equations Fail?)

Double shear layer example:



Why Does Filtering Work ? (Or, Why Do the Unfiltered Equations Fail?)

Consider the model problem:
$$\frac{\partial u}{\partial t} = -\mathbf{c} \cdot \nabla u$$

Weighted residual formulation: $B\frac{d\underline{u}}{dt} = -C\underline{u}$

$$B_{ij} = \int_{\Omega} \phi_i \phi_j \, dV = \text{symm. pos. def.}$$

$$C_{ij} = \int_{\Omega} \phi_i \mathbf{c} \cdot \nabla \phi_j \, dV$$

= $-\int_{\Omega} \phi_j \mathbf{c} \cdot \nabla \phi_i \, dV - \int_{\Omega} \phi_j \phi_j \nabla \cdot \mathbf{c} \, dV$
= skew symmetric_if $\nabla \cdot \mathbf{c} = 0$

 $B^{-1}C \longrightarrow$ imaginary eigenvalues

Discrete problem should never blow up.

Why Does Filtering Work ? (Or, Why Do the Unfiltered Equations Fail?)

Weighted residual formulation vs. spectral element method:

$$C_{ij} = (\phi_i, \mathbf{c} \cdot \nabla \phi_j) = -C_{ji}$$
$$\tilde{C}_{ij} = (\phi_i, \mathbf{c} \cdot \nabla \phi_j)_N \neq -\tilde{C}_{ji}$$

This suggests the use of over-integration (dealiasing) to ensure that skew-symmetry is retained

$$C_{ij} = (J\phi_i, (J\mathbf{c}) \cdot J\nabla\phi_j)_M$$

 $J_{pq} := h_q^N(\xi_p^M)$ interpolation matrix (1D, single element)

Aliased / Dealiased Eigenvalues: $u_t + \mathbf{c} \cdot \nabla u = 0$

- Velocity fields model first-order terms in expansion of straining and rotating flows.
 - For straining case, $\frac{d}{dt}|u|^2 \sim |\hat{u}_{N}|^2 + |\hat{u}_N|^2$
 - Rotational case is skew-symmetric.
 - Filtering attacks the leading-order unstable mode.



Stabilization Summary

- Filtering acts like well-tuned hyperviscosity
 - Attacks only the fine scale modes (that, numerically speaking, shouldn't have energy anyway...)
 - Can precisely identify which modes in the SE expansion to suppress (unlike differential filters)
 - Does not compromise spectral convergence
- Dealiasing of convection operator recommended for high Reynolds number applications to avoid spurious eigenvalues
 - Can run double shear-layer roll-up problem forever with

$$- v = 0$$
,

– no filtering

v = 0, no filter

 $V = 10^{-5}$, no filter

v = 0, filter = (.1,.025)







Linear Solvers

Navier-Stokes Solution Strategy

Semi-implicit: explicit treatment of nonlinear term.
Leads to Stokes saddle problem, which is algebraically split

MPR 90, Blair-Perot 93,

$$\begin{bmatrix} \mathbf{H} & -\mathbf{D}^T \\ -\mathbf{D} & 0 \end{bmatrix} \begin{pmatrix} \underline{\mathbf{u}}^n \\ \underline{p}^n - \underline{p}^{n-1} \end{pmatrix} = \begin{pmatrix} \mathbf{B}\underline{\mathbf{f}} + \mathbf{D}^T \underline{p}^{n-1} \\ \underline{f}_p \end{pmatrix}$$

$$\begin{bmatrix} \mathbf{H} & -\frac{\Delta t}{\beta_0} \mathbf{H} \mathbf{B}^{-1} \mathbf{D}^T \end{bmatrix} \left(\underline{\mathbf{u}}^n \\ \mathbf{D} & \begin{bmatrix} E \end{bmatrix} \left(\underline{p}^n - \underline{p}^{n-1} \right) = \left(\underline{g} \end{bmatrix} + \mathbf{D}^T \underline{p}^{n-1} \right) + \left(\underline{\mathbf{r}} \\ \underline{0} \end{bmatrix} \cdot ,$$

$$E := \frac{\Delta t}{\beta_0} \mathbf{D} \mathbf{B}^{-1} \mathbf{D}^T , \qquad \mathbf{\underline{r}} = O(\Delta t^2)$$

- **E** consistent Poisson operator for pressure, SPD
 - Stiffest substep in Navier-Stokes time advancement
 - Most compute-intensive phase

Couzv 95

- Spectrally equivalent to SEM Laplacian, A

Pressure Solution Strategy: $E\underline{p}^n = \underline{g}^n$

- 1. Projection: compute best approximation from previous time steps
 - Compute \underline{p}^* in span{ \underline{p}^{n-1} , \underline{p}^{n-2} , ..., \underline{p}^{n-l} } through straightforward projection.
 - Typically a 2-fold savings in Navier-Stokes solution time.
 - Cost: 1 (or 2) matvecs in E per timestep

2. Preconditioned CG or GMRES to solve

$$E \Delta \underline{p} = \underline{g}^n - E \underline{p}^*$$

Initial Guess for $A\underline{x}^n = \underline{b}^n$ via Projection (A-E, SPD)

$$\begin{aligned} \text{Given} \cdot \underline{b}^{n} & \cdot \{ \underline{\tilde{x}}_{1}, \dots, \underline{\tilde{x}}_{l} \} \text{ satisfying } \underline{\tilde{x}}_{i}^{T} A \underline{\tilde{x}}_{j} = \delta_{ij}, \\ \cdot \text{ Set } \underline{\tilde{x}} &:= \Sigma \alpha_{i} \underline{\tilde{x}}_{i}, \quad \alpha_{i} = \underline{\tilde{x}}_{i}^{T} \underline{b} \quad \text{(best fit solution)} \\ \cdot \text{ Set } \Delta \underline{b} &:= \underline{b}^{n} - A \underline{\bar{x}} \\ \cdot \text{ Solve } A \Delta \underline{x} &= \Delta \underline{b} \text{ to } tol \ \epsilon \quad \text{(black box solver)} \\ \cdot \underline{x}^{n} &:= \underline{\bar{x}} + \Delta \underline{x} \\ \cdot \text{ If } (l = l_{\max}) \text{ then} \\ \underline{\tilde{x}}_{1} &= \underline{x}^{n} / || \underline{x}^{n} ||_{A} \\ l &= 1 \\ \text{else} \\ \underline{\tilde{x}}_{l+1} &= (\Delta \underline{x} - \Sigma \beta_{i} \underline{\tilde{x}}_{i}) / (\Delta \underline{x}^{T} A \Delta \underline{x} - \Sigma \beta_{i}^{2})^{\frac{1}{2}}, \quad \beta_{i} &= \underline{\tilde{x}}_{i} A \Delta \underline{x} \\ l &= l+1 \end{aligned}$$

CHUII

Initial guess for Epⁿ = gⁿ via projection onto previous solutions ICASE TR '93; CMAME '98

 $\blacksquare \| \underline{p}^n - \underline{p}^* \|_E = O(\Delta t^l) + O(\varepsilon_{\text{tol}})$

two additional mat-vecs per step

storage: $2+l_{max}$ vectors

results with/without projection (1.6 million pressure nodes)



• 4 fold reduction in iteration count, 2-4 in typical applications

Overlapping Additive Schwarz Preconditioner / Smoother

⁽Dryja & Widlund 87, Pahl 93, PF 97, FMT 00)



Local Overlapping Solves: FEM-based Poisson problems with homogeneous Dirichlet boundary conditions, A_e .

Coarse Grid Solve: Poisson problem using linear finite elements on entire spectral element mesh, A_0 (GLOBAL).

E-Based Schwarz vs. SEMG for High-Aspect Ratio



Solver Performance: Hybrid Schwarz-Multigrid

Magneto-rotational instability

(Obabko, Cattaneo & F.)

- E=140000, N=9 (n = 112 M), P=32768 (BG/L)
- ~ 1.2 sec/step
- ~ 8 iterations / step for U & B
- Key is to have a scalable coarse-grid solver





Navier-Stokes Solution time break down for n=120 M, $n_c = 417000$

Case/P	Total	QQ^T	Coarse	all_reduce()
x4096	1994	125	1180	1.2
a4096	1112	125	192	1.4
b4096	846	126	25	1.
8192	460	88	22	1.
16384	266	64	20	1.

Cannot consider XX^T on larger problems.

"a4096" case is relies on pairwise + all_reduce

 First version (not shown, pairwise-only), was not much faster than XX^T. Why?

Outline

Nek5000 capabilities

Equations, timestepping, and SEM basics

Parallel / serial issues that you should understand

Serial / Parallel Issues

Locally, the SEM is **structured**.

Globally, the SEM is **unstructured**.

Vectorization and serial performance derive from the structured aspects of the computation.

Parallelism and geometric flexibility derive from the unstructured, element-by-element, operator evaluation.

Elements, or groups of elements are distributed across processors, but an element is never subdivided.

Parallel Structure

Elements are assigned in ascending order to each processor



Parallel, local element numbering

Parallel Structure



For the most part, don't care about global element *numbering*

- (We'll show some examples where one might)

Key point is that,

- on proc 0, nelt=2 (nelt = # elements in temperature domain)
 - on proc 1, nelt=3 (nelv = # elements in fluid domain, usually = nelt)

Parallel Structure



Arrays that distinguish which processor has which elements:

- proc 0	proc 1
• nelt=2	nelt=3
• lglel=(2,5)	lglel=(1,3,4)

Common arrays (scaling as nelgt, but only two such arrays):

- gllel=(1,1,2,3,2), gllnid=(1,0,1,1,0)

Serial Structure

All data contiguously packed (and quad-aligned):

real u(lx1,ly1,lz1,lelt)

 Indicates that u is a collection of elements, e=1,...,Nelt =< left, each of size (N+1)^d, d=2 or 3

Serial / Parallel Usage

A common operation (1st way...)
Parallel Version

s=0 do e=1,nelv do iz=1,nz1 do iy=1,ny1 do ix=1,nx1 s=s+u(ix,iy,iz,e)enddo,...,enddo

s=0do e=1,nelv
do iz=1,nz1
do iy=1,ny1
do ix=1,nx1 s=s+u(ix,iy,iz,e)enddo,...,enddo

s=glsum(s,1)

Serial / Parallel Usage

A common operation (2nd way...)
Parallel Version

n=nx1*ny1*nz1*nelv s=0 do i=1,n s=s+u(i,1,1,1)enddo n=nx1*ny1*nz1*nelvs=0 do i=1,n s=s+u(i,1,1,1)enddo

s=glmax(s,1)

Serial / Parallel Usage

A common operation (3rd way...)
Parallel Version

n=nx1*ny1*nz1*nelv

n=nx1*ny1*nz1*nelv

s=glsum(u,n)

s=glsum(u,n)

- If you want a *local* max:

s=vlsum(u,n)

– Note: Important that every processor calls glmax()!!

Considerations as you move forward

For performance:

- Optimal tuning for mxm ?
- What about the pressure solve?
 - Projection on?
- For stability:
 - Dealiasing on?
 - Filtering on?
- Open issues we expect to address in the next 12-24 months:
 - SEMG for the P_N - P_N formulation
 - (near) monotone advection schemes

Thank You!